SKOBELKIN, O.K., kand.med. nauk (Kalinin)

Surgical complications in ascariasis. Med. sestra 22 no.10t
(MIRA 16:12)

SKOBKLKIN, O.K. (Kalinin, 2-ya ul. Shevchenko, d.40, kv.25)

Extensive traumatic diaphragmatic hernia in a child. Vestm. khir. Grekov. 90 no.4291 Ap163 (MIRA 17:2)

1. Iz fakul tetskoy khirurgicheskoy kliniki (zav. - prof. A.G. Karavanov) Kalininskogo meditsinskogo instituta na baze Kalininskoy oblastnoy bol nitsy No.1 (glavnyy vrach - zasluzhennyy vrach RSFSR A.A. Sokolov).

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

SKOBELKIN, 5.K., kand. med. nauk (Kalinin, 5, 2-ya ulitsa Shevehenko, d.40, kv.25)

**电视电视电视** 

Choledochdodenal fistulas. Vest. Khir. 91 no.10:113-114 0 163. (MIRA 17:7)

1. Iz fakul'tetskoy khirurgicheskoy kliniki (ispolnyayushchiy obyazannosti zaveduyushchego - dotsent N.V. Zavadovskaya) Kalininskogo meditsinskogo instituta (rektor - dotsent A.N. Kushnev) na baze oblastnoy klinicheskoy bol'nitay (glavnyy vrach - zasluzhennyy vrach RSFSR A.A. Sokolov).

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

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SKOBELKIN, O.K., kand.med.nauk

Revascularization of a mobilized stomach with the aid of the splenic artery. Trudy KGMI no.10:372-374 163. (MIRA 18:1

1. Iz kafedry faku! 'te skoy khirurgii (zav. kafedroy zasluzhennyy deyatel nauki RSFSR - prof. V.S.Semenov) i kafedry torakalnoy khirurgii i anesteziologii Ukrainskogo tsentral nogo instituta usovershenstvovaniya vrachey (zav. kafedroy - prof. A.A.Shalimov).

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

L 20270-65 ACCESSION NR: ARLIQUESOS \$\0599\@\\\000\old\\\\051\\\\\051\\\\\\\

SOURCE: Ref. zh. Biologiya. Svodnyay tom, Abs. 1410.58

5 13

AUTHOR: Skobelkin, O. K.

TITLE: Circular defect esophagoplasty with a stomach saction on a

vascular pedicle

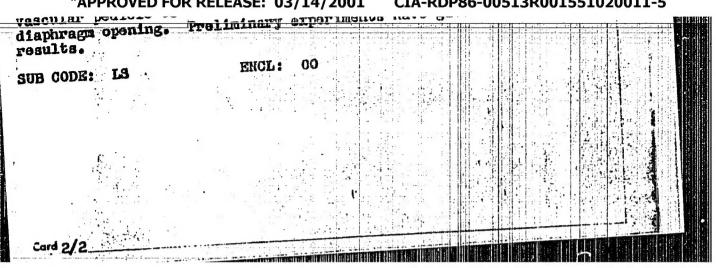
or mov, 1963. Yerevan, 1963, 375-378

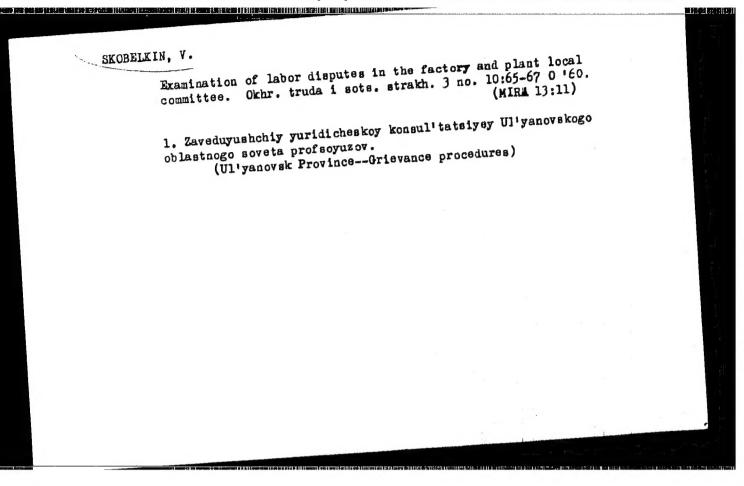
cadaver, stomach

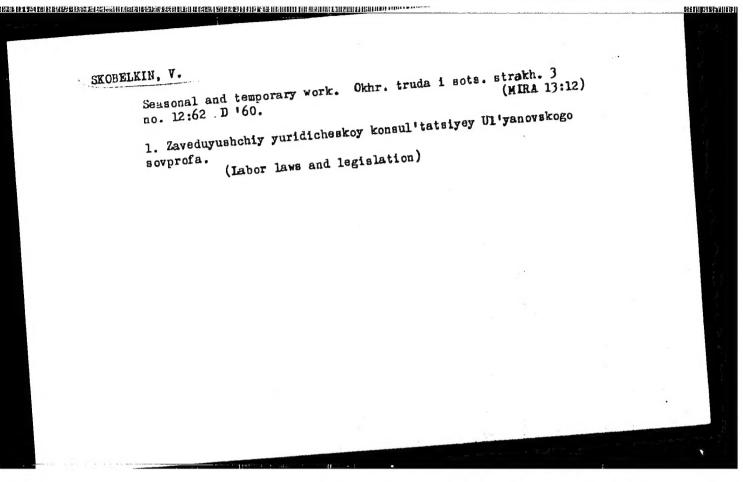
TRANSLATION: In experiments on 10 dog cadavers and 5 live dogs it was domonstrated that it is more effective to make a stomach tube resection from the large curvature of the atomach because of its greater length and size and highly adequate number of vessels. After mobilization and spleen removal, the vascular archie make it possible to bring the transplant into the thoracic cavity to the level of the arch of the aorta and higher. The distal end of the transplant

L 20270-65
ACCESSION MR: ARHOUS869
started 2 to 3 om higher than the pylorus and the proximal and started 2 to 7 cm long started 2 to 3 cm higher than the pylorus and the proximal and long started 2 to 7 cm long started 2 to 7 cm long

### CIA-RDP86-00513R001551020011-5 "APPROVED FOR RELEASE: 03/14/2001



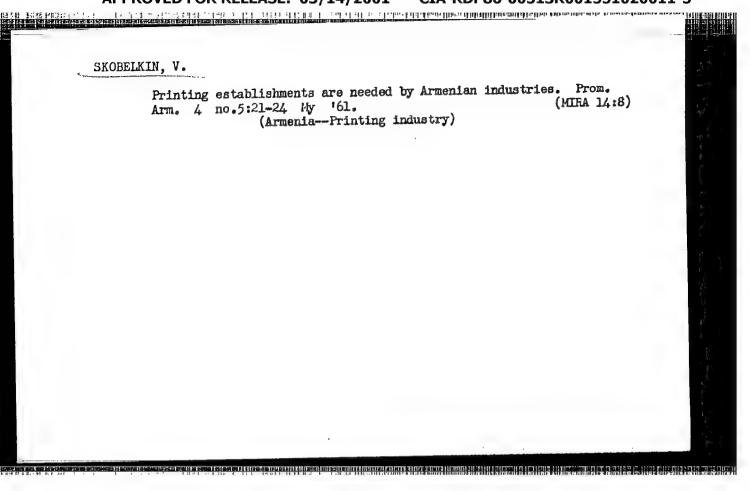




GERSHANOV, Ye.; SKOBELKIN, V.

Commission for Labor Disputes at enterprised and institutions. Sots.
trud 5 no.1:140-145 Ja '60. (MIRa 13:6)

(Grievance procedures)



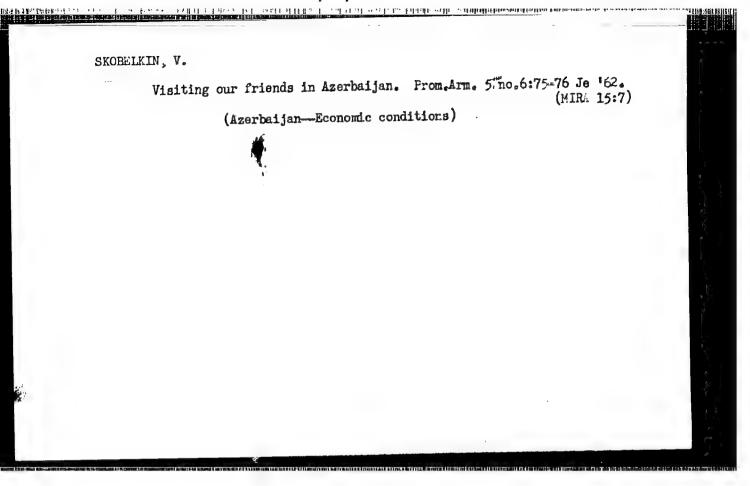
SKOBELKIN, V.

Consulation by the Ul'yanov Province Trade-Union Council. Okhr. truda i sots. strakh. 4 no.5:53-54 My '61. (MIRA 14:5)

1. Zaveduyushchiy yuridicheskoy konsul'tatsiyey Ul'yanovskogo oblsovprofa. (Labor laws and legislation)

Is the court right? Okhr.truda i sots.strakh. 5 no.4142-43
Ap '62.

1. Zaveduyushchiy yuridicheskoy konsul'tatsiyey Ul'yanovskogo oblsovprofa.
(Maternal and infant welfare) (Employees, Dismissal of)



### SKOBELKIN, V.

Liability responsibility of the person reponsible for an industrial acciedent and illegal dismissal. Okhr.truda i sots strakh. 5 no.10:42 0 '62. (MIRA 15:11)

1. Zaveduyushchiy yuridicheskoy konsul'tatsiyey Ul'yanovskogo oblastnogo soveta professional'nykh soyuzov.

(Employers' liability) (Employees, Dismissal of)

SKOBELKIN, V., kand.yuridicheskikh nauk

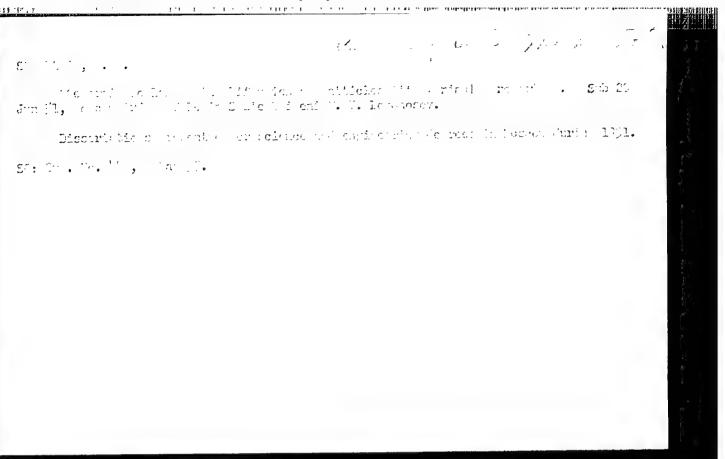
Procedure for bringing a damage suit against a defendant who caused an industrial accident. Okhr. truda i sots. strakh. 6 no.9:43-45 (MIRA 16:10)

SKOBELKIN, V. I.

and A.G. Shafigullin- authors of a-Review & Summary of "Theory of Chain Processes" by N.S. Akulov; State Pub. House of Tech & Theoret. Lit., Moscow, 1951.

Zhur Fiz Khim, Vol XXVII, #1, pp 151-2

W-30868, 18 Aug 54



SKOBELKIN, V. I. Sr. Sci. Colleague

"Some Questions of Non-Linear Electrodynamics," a paper given at the All-University Scientific Conference "Lomonosov Lectures", Vest. Mosk. Un., No.8, 1953.

Trabslation U07895, 1 Mar 56

- 1. SMOBELHIN, V. I., SHAFIGULIIN, A. G.
- 2. USSR (600)
- 4. Akulov, Nikolai Sergeevich
- N. S. Akulov's book "Theory of chain reactions." Reviewed by V. I. Skobelkin, A. G. Shafigullin. Zhur. fiz. khim. 27, no. 1, 1953.

9. Monthly List of Russian Accessions, Library of Congress, May 1953, Unclassified.

and the fire and the first of the model of the fire of

USSR/Physics - Vacuum pump

FD-1C83

Card 1/1

Pub. 153 - 19/24

Author

: Skobelkin, V. I., and Yushchenkova, N. I.

Title

: Theory of the vapor-jet vacuum pump

Periodical

: Zhur. tekh. fiz., 24, No 10, 1879-1891, Oct 1954

Abstract

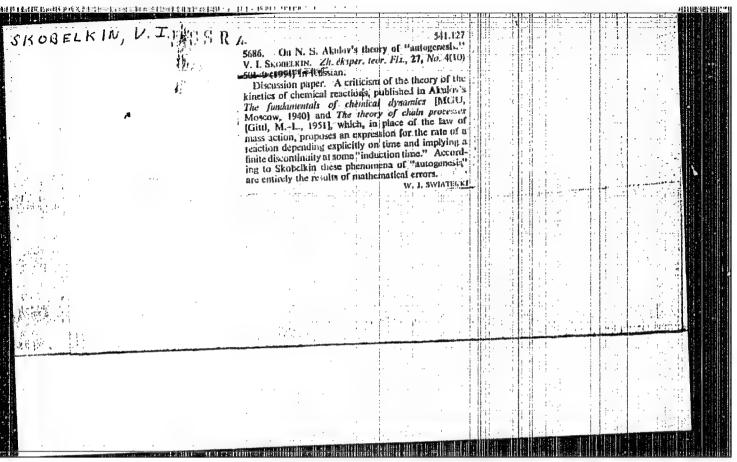
: The authors investigate the interaction between the gas to be pumped out and the supersonic vapor jet. They clarify the mechanism governing the process and thus are enabled to calculate the speed of pumping out of the gas and to determine the influence of the various parameters upon this speed. They note that their results differ from those obtained by

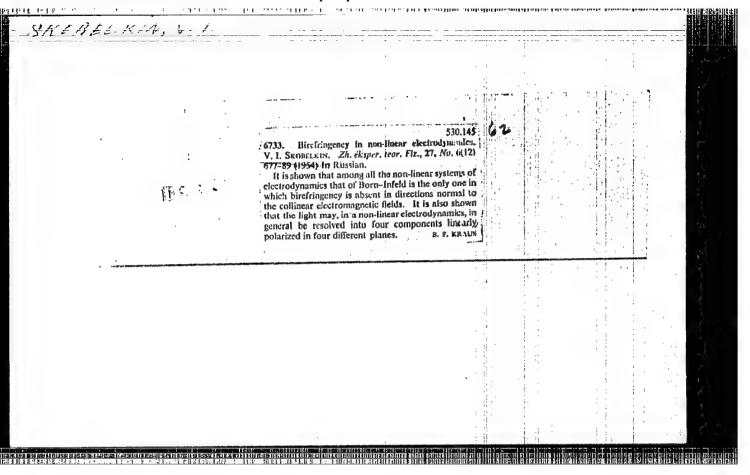
the USSR authors Lifshits and Rozentsveyg (ibid., No 8, 1952).

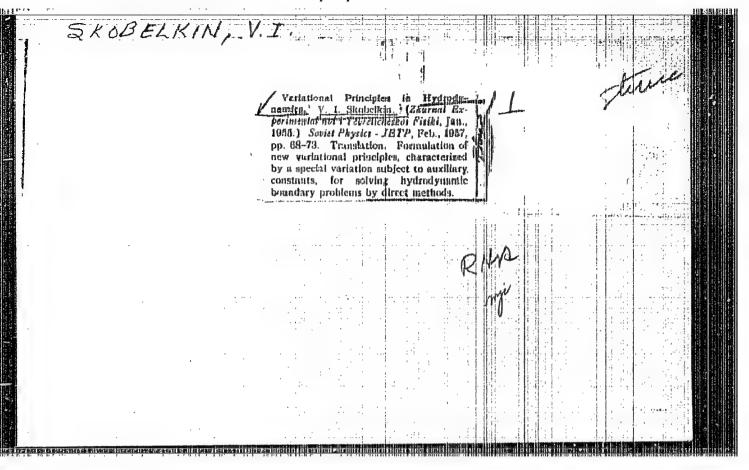
Institution : -

Submitted : April 3, 1953

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"







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Category : USSR/Atomic and Molecular Physics - Gases

D-7

Abs Jour : Ref Zhur - Fizika, No 1, 1957 No 932

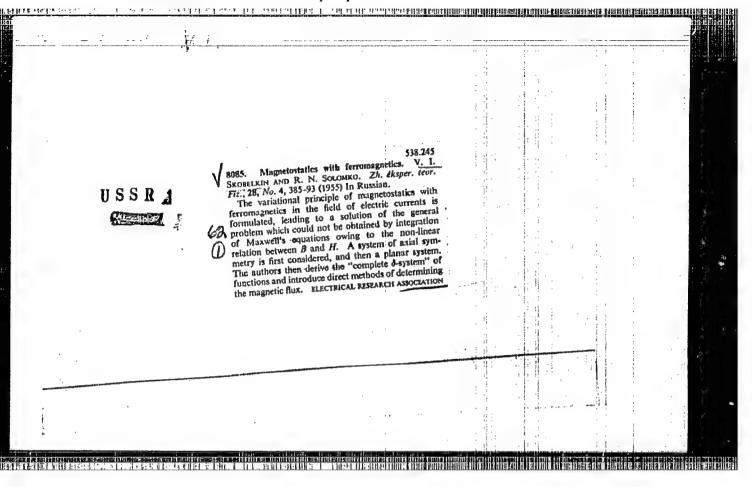
Author : Skobelkin, V.I., Yushchenkova, N.I.

Title : Corrections to Article "Theory of Vapor-Jet Vacuum Pump."

Orig Pub : Zh. tekhn. fiziki, 1955, 25, No 2, 66

Abstract : Refers to Ref. Zhur. Fiz. 1955, 8952

Card : 1/1



CARD 1 / 2

SUBJECT USSR / PHYSICS

PA - 1537

AUTHOR SKOBELKIN, V.I.

On Various Principles in Hydrodynamics

TITLE On Various Principles in hydrodynamics (1956) Zurn.eksp.i teor.fis, 31, fasc.2, 317-323 (1956)

Issued: 5.10.1956

In the works by several authors the variation functional contained some field parameters which were then varied independently (density, velocity, pressure, and others, eg. density and stream functions). However, the number of parameters used on this occasion is not minimal. It is shown that for the construction of the functional in the general case of a threedimensional motion two defining quantities of the field (stream-functions) suffice.

The principle of the lowest flow potential:  $\overline{Y}$ ,  $\overline{Y}$  are assumed to be the parameters for the determination of a certain streamline of a steady hydrodynamic field. An orthogonal trihedron  $\overline{x}_i$  is constructed at each point of the streamline, and to each trihedron an energy-momentum-tensor  $T_{ik}$  is assigned. The quantities  $\overline{Y}(x_i)$ ,  $\overline{X}(x_i)$  of the field satisfy the relations  $\overline{Y}(x_i)$ ,  $\overline{X}(x_i)$  of the field satisfy the relations  $\overline{Y}(x_i)$ ,  $\overline{Y}(x_i)$  of the mechanic energy in a real steady field along the streamline it is then true that:  $\int_{0}^{\infty} T_{ik} / \partial \overline{x}_k dS = (\overline{Y}^2/2) + (dP/Q) + R(Q, \overline{Y}, \overline{X}) = E(\overline{Y}, \overline{X})$ . Here R is the work of friction forces per unit of measure of the liquid. The LAGRANGIAN is  $\overline{Y}_{ik} = \overline{Y}_{ik} = \overline{Y}_{$ 

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Zurn.eksp.i teor.fis, 31, fasc.2, 317-323 (1956) CARD 2 / 2 PA - 1537 field. Herefrom follows:  $\frac{\partial T_{SS}}{\partial \Psi} - (\partial/\partial x_i \partial) \left(\frac{\partial T_{SS}}{\partial \Psi} - (\partial/\partial x_i \partial) (\frac{\partial T_{SS}}{\partial \Psi} - (\partial/\partial x_i \partial) (\frac{\partial$ 

INSTITUTION: Moscow State University.

PA - 1218 CARD 1 / 2 USSR / PHYSICS SUBJECT SKOBELKIN. V.I.

AUTHOR The Principle of the Smallest Flow Potential. TITLE

Dokl. Akad. Nauk, 108, 787-790 (1956) PERIODICAL reviewed 8 / 1956 Publ. 6 / 1956

The present work establishes and formulates a variation principle for the solution of various boundary value problems of gas dynamics by means of new and partly direct methods. The parameters T and A are assumed to define a certain streamline in a steady flow of a perfect gas. The surfaces  $\Psi = \text{const}$  and  $\mathcal{N} = \text{const}$  form the flow surfaces.  $\Psi$  and  $\mathcal{N}$  are chosen in such a manner that  $\varrho V_{\chi} = D(\Psi, \mathcal{N})/D(y,z)$ ,  $\varrho V_{\chi} = D(\Psi, \mathcal{N})/D(x,y)$ .

The couple of functions  $(\Psi, \stackrel{\frown}{\checkmark})$  may be considered to be the flow functions in a threedimensional motion. The continuity equation div  $Q \stackrel{\frown}{V} = 0$  is satisfied identically for any  $\nabla$ ,  $\hat{\mathcal{N}}$ . If on a boundary surface  $\sigma$  the distribution of the flow (the distribution of Y and ) and the distribution of the total mechanical energy  $\mathcal{E}(\Psi, \mathcal{C})$  referred to units of measure of the gas are known, the

BERNOULLI equation has the form  $(v^2/2) + \int dP/Q = \mathcal{E}(\Psi, \mathcal{F})$ . Here P is the gas pressure, and integration is carried out along the streamline. Furthermore, it applies for the energy theorem (1.principal theorem) that dQ = dU + P d(1/Q). Here Q denotes the quantity of heat and U the inner energy.

If  $\Psi$  and  $\vartheta$  are varied within the domain  $\Omega$  (in which the motion of the gas

CARD 2 / 2 PA - 1218 Dokl. Akad. Nauk, 108, 787-790 (1956) has no strong or slight discontinuities) while their values are maintained on the surface  $\sigma$  the following principle of variation applies for the motion of the gas in  $\mathcal{H}$ : The integral  $I = \int_{\Omega} (P + \varrho V^2) d \omega$ , which is equal to the work performed by the entire impulse flow of the directioned motion of gas (flow potential), assumes a steady value in the case of an actual motion of a perfect gas which has become steady. We now introduce the geometrical characteristic of flow:  $\theta = (D(\Psi, \mathcal{S})/D(y,z))^2 + (D(\Psi, \mathcal{S})/D(z,x))^2 + (D(\Psi, \mathcal{S})/D(x,y))^2 \text{ and restrict our}$ attention to the motion which is barotropic for a given gas particle. The OSTROGRADSKIJ equations corresponding to the equation  $\delta I = 0$  are given, several times transformed, and specialized for an adiabatic motion. It is shown that, in the case of adiabatic subsonic motions of the gas (which are usually vortex-like), the solution of the equations represents a sharp minimum of the functional I. Proof bases upon WEIERSTRASS function of the calculus of variations, and is followed step by step. In the case of a subsonic motion the phase surface Z=L (0, Y, N) is convex, but in the case of supersonic motions without strong or slight discontinuities the phase surface has a saddle at all of its points. If boundary conditions for Y and con the surface of are given, there is only one solution.

INSTITUTION: Moscow State University "M.V.LOMONOSOV"

PREPERA A F L.

124-11-12524D

Translation from: Referativnyy Zhurnal, Mekhanika, 1957, Nr 11, p 31 (USSR)

AUTHOR: Skobelkin, V. I.

TITLE: The Principle of the Minimal Stream Potential and its Application to

Problems in Gasdynamics, Electrodynamics, and Combustion.

(Printsip naimen'shego potentsiala toka i yego prilozheniya k zadacham

gazovoy dinamiki, elektrodinamiki i goreniya)

ABSTRACT: Bibliographic entry of the Author's dissertation for the degree of

Doctor of Physical-Mathematical Sciences, Institute of Mechanics,

Academy of Sciences, USSR, Moscow, 1957.

ASSOCIATION: Institute of Mechanics, Academy of Sciences, USSR, Moscow

2 Card 1/1

SKCHLER, V.I., Dec Physel th Sci-(disc) "Variation principle in hydrodynchics." los, 1950. 22 pp (Acad Sci USSR. Inst of Lachanics), no coice. Bibliography: 59 21-22 (29 titles) (KL, 26-50,104)

AUTHORS:

Kogarko, S. M., Skobelkin, V. I.

SOV/20-120-6-32/59

TITLE:

Relaxation Interaction Between Shock Waves and the Combustica Zone (Relaksatsionnoye vzaimodeystviye udarnykh voln s zone)

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PERIODICAL:

Doklady Akademii nauk SSSR, 1958, Vol 120, Nr 6,

pp 1280 - 1283 (USSR)

ABSTRACT:

The paper under review presents an investigation of the influence of the kinetics of the combustion upon the structure and the intensity of a shock wave when it passes through the combustion zone. The relaxation time  $\mathbf{t_r}$  (the period during

which no noticeable influence is exerted upon the diffusion currents and the heat currents in the reaction zone by the reaction conditions suddenly modified by the shock wave) is of the same order as the reaction period  $\mathcal{T}$  (10<sup>-3</sup> - 10<sup>-5</sup> sec).  $\mathcal{T}$  is defined as the ratio of the width of the combustion zone  $\ell$  and the normal expansion velocity of the flame. During  $t_r$ 

the temperature and the pressure within the reaction zone increase. During the relaxation time the excess momentum in the

Card 1/3

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Relaxation Interaction Between Shock Waves and the Combustion Zone

SOV/20-120-6-32/59

reaction zone is transformed into a shock wave propagating forward and backward from the reaction zone. Equations giving the state in the shock wave are written down. The whole amount of heat liberated in the passage of the shock wave is expended for the increase of the internal energy of the gas in the reaction zone.  $\mathcal{T}_{\mathbf{h}}$  denotes the duration of the passage of the shock wave through the combustion zone. If  $\tau > \tau$  , the reaction is not completed during the passage of the wave through the front of the flame and only a certain proportion of the chemical energy which is expended for the increase of the momentum of the wave is imported to the wave. If  $\tau \leqslant \tau$  , the reaction is completed within the period  $\tau$  . The total momentum of the shock wave after passing the combustion zone is combined from the interest momentum I and the relaxation momentum I. The maximum amplification of the momentum of the shock wave at T =T, may be termed momentum resonance. Finally a method for the determination of the index of refraction is presented.

Card 2/3

Relaxation Interaction Between Shock Waves and the Sombustion Zone

SOV/20-120-6-32/59

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There is 1 figure.

ASSOCIATION:

Institut khimicheskoy fiziki Akademii nauk SSSR (Institute of

Chemical Physics, AS USSR)

PRESENTED:

March 6, 1958, by N. N. Semenov, Member, Academy of Sciences,

USSR

SUBMITTED:

February 25, 1958

1. Shock waves--Analysis 2. Combustion--Analysis 3. Mathematic 18

---Applications

Card 3/3

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

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5(4), 24(8)AUTHOR:

Skobelkin, V. I.

SOV/20-122-3-30/57

TITLE:

On the Thermodynamic Equilibrium of the Surfaces of a Strong Explosion (O termodinamicheskom ravnovesii poverkhnostey

sil'nogo rasryva)

ումին ու ու արագարան անանակարան արագարանի արանի արանական արանական անական հայարական հայարական հայարական հայարակ

PERIODICAL:

Doklady Akademii nauk SSSR, 1958, Vol 122, Nr 3, pp 431-433

(USSR)

ABSTRACT:

In hydrodynamics, the problem of the thermodynamic equilibrium appears as often as a variation principle is formulated for the description of the gas motion. The infinitely small deflections from the actual motion for given boundary and initial conditions must satisfy not only the variation equation which is equivalent to the equation of motion but also to the variation equation  $\Delta Q - T\delta S = 0$ . This equation is equivalent to the second law of thermodynamics for a real state of a gas in thermodynamic equilibrium. The above-given equation, therefore, limits the class of the permissible variations. This paper deals with the general case in which the gas moves in the presence of a shock wave or of a sharply defined front

Card 1/3

of reaction (which divides the initial gas from the burnt

SOV/20-122-3-30/57
On the Thermodynamic Equilibrium of the Surfaces of a Strong Explosion

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gas). The laws for the conservation of mass, momentum, and energy are given explicitly and then the Hugoniot (Gyugonio) equation is derived from them and a variation equation is given for the real subsonic flow of an ideal gas. To this variation equation, the initially given variation equation  $\Delta Q$  - ToS has to be added. As an example, the author investigates the subsonic flow of a perfect gas in a tube in the presence of a flame front which propagates with a given velocity in an isentropic gas with a given Poisson (Puasson) constant. A figure shows the Hugoniot adiabates for the shock wave, for the detonation wave, and for a weak deflagration. According to the deliberations of this paper, the model of the notion of 2 incompressible liquids does not satisfy the equation of the thermodynamic equilibrium of the flame front. If, however, the gas moves in the presence of a shock wave, the non-equilibrium conditions in a transition layer of the order of the free length of path of the gas molecules has to be investigated for the determination of  $\Delta$  Q of an elementary particle behind the shock wave. There are 1 figure and 12 references, 7 of which are Soviet.

Card 2/3

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\$507/20-122-3-30/57\$ On the Thermodynamic Equilibrium of the Surfaces of a Strong Explosion

ASSOCIATION: Moskovskiy gosudarstvennyy universitet im. M. V. Lomonosova

(Moscow State University imeni M. V. Lomonosov)

PRESENTED: June 30, 1958, by S. A. Yekshinskiy, Academician

SUBMITTED: May 26, 1958

Card 3/3

507/20-122-6-25/49 5(4), 10(7)Kogarko, S. M., Skobelkin, V. I., Kazakov, A. N. AUTHORS:

The Interaction Between Shock Waves and the Front of a Flame TITLE:

(Vzaimodeystviye udarnykh voln s frontom plameni)

Doklady Akademii nauk SSSR, 1958, Vol 122, Nr 6, pp 1046-1048 PERIODICAL:

(USSR)

The present paper investigates the intensification of shock ABSTRACT:

waves in their interaction with the front of a flame by

variation of the normal combustion process in the shock wave. The length of the shock wave is assumed to be sufficient in the direction of the reaction zone. For the interaction between such a shock wave and the flame front the following applies: 1) The shock wave is transformed at the flame front (like on the boundary dividing two media). In this way a refracted and a reflected wave are formed. The flame front can by approximation be considered to be a contact-discontinuity. The expressions for the refraction coefficient are written down. 2) When passing through the flame front the shock wave

compresses the gas in the reaction zone, whereby temperature rises. This temperature rise increases reaction velocity, so

Card 1/3

507/20-122-6-25/49

The Interaction Between Shock Waves and the Front of a Flame

that the propagation velocity of the flame is also increased. This propagation velocity increases very rapidly, and therefore this process may be looked upon as a sort of explosion in the gas current behind the shock wave; it causes the formation of 2 additional (intensifying) shock waves. The shock wave front moves with subscnic velocity in relation to the disturbed gas, and therefore any kind of disturbance is able to catch up with this front in the current behind the shock front, thus changing its structure. The propagation velocity of the flame is not increased immediately upon arrival of the shock wave, but only after a certain relaxation time. The latter is of the same order of magnitude as the duration of reaction. A diagram schematically shows the intensification of the shock wave when passing through the flame front. Expressions for shock front calculation are given. The new propagation velocity of the flame is calculated according to the theory developed by Zel'dovich. The amplitude of the intensifying shock wave depends upon the amplitude of the initial shock wave as well as on the kinetic properties (reaction velocity, calorific value, activation energy, etc.) of the fuel. The second diagram shows the amplitude of the inten-

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507/20-122-6-25/49

The Interaction Between Shock Waves and the Front of a Flame

sifying shock wave of compression in the reaction zone for 2 different propagation velocities. There are 2 figures and

5 Soviet references.

ASSOCIATION: Institut khimicheskoy fiziki Akademii nauk SSSR

(Institute for Chemical Physics of the Academy of Sciences,

USSR)

PRESENTED: June 21, 1958, by V. N. Kondrat'yev, Academician

SUBMITTED: June 11, 1958

Card 3/3

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	MARK X 3000	deregentally po megationy gravestatata, side, 250., (Special Plancy; truly Konferentelli, [Problem & Negation/checker and Pasan Dynamic; Trussactions of a Gantlemee) Maga, Independent and Pasan Spinalis; Trussactions of a Exruta ally laserted, 1,000 copies printed.	Honsoring Agency: Akademiya nank Latviyelory SER, Testitut fisiki,	Mitorial beard: D.A. Frank-Kammetskiy, Doctor of Physics and Matheastics, Professor; A.L. Vol'dat, Doctor of Jechnical Sciences, Professor; I.M. Kirto Border of Physics and Mathematics; 7.7%, Welline, Candidate of Physics and Mathematics; VG, Vitol, Gandidate of Physics and Matheastics; Publ. Kruntin and V.Is. Krevolesno.	MA. 1 A. Poytal Pouns Doch, Ed. 1 A. Alyuvinya	FURINGEST This book is intended for paysiciats working in the field of magneto-daylord and places between the state of a conference beld in Rigs, June 1996, on problems in spyled est theoretical magnetobedynatics. The shates 1996 and theoretical magnetobedynatics. The shates 1996 and theoretical magnetobedynamics. The shates is the conference over the investigation of the basic tracks in theoretical places in the last tracks in the restate and spreading dieg research in different bravels of the conference of the restaint bravels and presenting the participation of theoretical places from the restaint to the problem of the problem of the restaint in the special places from the restaint of the forter thio conference as the restaint of the forter thio says to be half a requirely in the fitter; for any man, since the forter the last sequence is not payers and conference as the branch of the first part deals with propers and commune on payers are presented by the embory themselves in an abridged form. The book is divided into the problem in the first part deals with problem in theoretical angelves problem as the application of the fraction of the problem in the special magneto-dredynamics and plants and plants and the first part deals with problem in the laporate and angelves/dredynamics of plants in a series of the appearance of the special of the fractions of the problem in the series of the fractions of the dress of the sealers of the series and the state of the sealers of the search of the sealers of the searth of the sealers of the sealers of the sealers of the sealers	Syronstally, S.I. On the Stability of Shock Fewe is Magnetiahydrodynandes	Polerin, N.V., and G.To. Lybership, The Impossibility of Rave- Satisfiathon Baves in Magaschyddynaufes	lordanskip, 6.V. Langlan's Theorem in Magnetohydrodynamics	Eissige, Mile, and Vit. Insplymer. Chique Shork Were in Planna Vith Finite Combetterry and Certain Other Problems in Shock-Were Theory	Stobelling Late. Thermodynamic Equilibrium of Strong Discontinuity Furtaces	Abbiyees, A.I., and A.G. Sitenico. On the Theory of Bydrouagnetic	Morosov, Asf. Cherenkov Generation of Magnetic found Maves	#Sterko, A.O., and Yu.A. Kirothin, Disaljation of Epiromagnetic Waves is furbulent Fluctuations	Cart 6/12	

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24(3) AUTHOR:

Skobelkin, V. I.

SOV/20-128-2-16/59

TITLE:

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On the Function of Magnetic Flux in a Three-dimensional Field

PERIODICAL:

Doklady Akademii nauk SSSR, 1959, Vol 128, Nr 2, pp 280-283

(USSR)

ABSTRACT:

In an article by V. I. Skobelkin and R. N. Solomko (Ref 1) a principle of variation wasformulated according to which a method was devised which permits an approximate solution of boundary problems of two-dimensional and axially symmetrical ferromagnetic systems within the current field. In both cases, the field was determined by one single function of the magnetic flux  $\Phi$  (which is similar to the method employed in hydrodynamics). In the case of two-dimensional ranges, in which magnetic permeability is constant and no currents are present, the function of magnetic flux may be regarded as the imaginary part of the complex potential (which satisfies the conditions of Riemann-Cauchy). In the general case, two everywhere steady functions  $\psi$  and  $\theta$  of the magnetic flux are to be introduced for a description of the three-dimensional field. The latter are connected with the magnetic induction  $\Phi$  by the relation

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On the Function of Magnetic Flux in a Three-dimensional Field

 $\vec{B} = \nabla \psi \times \nabla \vartheta$ . The planes  $\psi = \text{const}$  and  $\vartheta = \text{const}$  are the planes of magnetic flux. The intersections of any two planes  $\psi$  and  $\vartheta$  form force lines in the field. With the help of the functions  $\psi$  and  $\vartheta$  it is possible to plot a Lagrangian for the three-dimensional field, and the variation principle formulated in the afore-mentioned previous article can be explained for the general case. For the plotting of the Lagrangian the "vector of magnetic flux"  $\vec{R} = (1/2c)(\psi \nabla \vartheta - \vartheta \nabla \psi)$  and the "density of the potential function of the currents"  $u = -j\vec{R}$  are introduced, where  $\vec{J}$  denotes the vector of current density. If the density of magnetic energy of the field is expressed by

 $w = \int_0^{B^2} \overrightarrow{H} \ d\overrightarrow{B}, \ L = w + u \ holds \ for the Lagrangian. The variation principle may be formulated in the following manner: Among all solenoidal fields of magnetic induction possible (when the given boundary conditions for the real field of the closed magnetostatic system are satisfied), the space integral <math display="block">E = \int_{\Omega} L \ d \ \omega \ assumes \ the \ least \ value. The mathematical formulation of this equation is reduced to the equation <math>\delta E = 0$  under

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SOV/20-128-2-16/59

On the Function of Magnetic Flux in a Three-dimensional Field

the condition div  $\overrightarrow{B} = 0$ , which is satisfied in a similar way if  $\psi$  and  $\mathcal{A}$  are chosen as the field-determinant quantities. L may be expressed by  $(B^2 dB^2 + \sqrt{3} + \sqrt{3}) = 0$ 

be expressed by  $L = -\frac{1}{2c} \vec{J}(\psi \nabla \cdot \hat{y} - \hat{y} \nabla \cdot \psi) + \frac{1}{4\pi} \int_{0}^{B^{2}} \frac{dB^{2}}{2\mu(B^{2})} = L(\vec{J}, \psi, \hat{y}, \nabla, \psi, \nabla, \hat{y}), \text{ and it}$ 

is possible to choose such a plane of that  $n \times H = 0$  holds. n denotes the normal on those planes on which the magnetic permeability  $\mu$  or the current density j are unsteady. As a result, elliptical equations are obtained for the determination of  $\psi$  and  $\theta$  as well as the pertinent natural boundary conditions. The above equations for  $\psi$  and  $\theta$  determine the projections of the curl of magnetic field strength to the lines which are perpendicular to the surface of magnetic flux. To prove the existence of a distinct minimum of E and of the exclusion principle for the field strength  $\overline{H}$ , the Weierstrass function E of the variation problem must be constructed. The article is concluded with a discussion of the required conditions. There are 6 references, 5 of which are Soviet.

Card 3/4

30V/20-128-2-16/59 On the Function of Magnetic Flux in a Three-dimensional Field

Moskovskij gosudarstvennyy universitet im. M. V. Lomonosova ASSOCIATION:

ACHICLE IN THE PROPERTY OF THE

(Moscow State University imeni M. V. Lomonosov)

May 18, 1959, by S. A. Vekshinskiy, Academician PRESENTED:

May 18, 1959 SUBMITTED:

Card 4/4

SOV/20-128-3-22/58 24(5) Skobelkin, V. I. AUTHOR: Propagation of Vector Waves in Nonlinear Mesodynamics TITLE: Doklady Akademii nauk SSSR, 1959, Vol 128, Nr 3, pp 514-51-PERIODICAL: (USSR) The nonlinear mesodynamics of the scalar and pseudoscalar ABSTRACT: field was investigated in an article by D. I. Blokhintser (Bo The generalization of linear vectorial mesodynamics consists in an investigation of the arbitrary function  $\begin{cases} \begin{cases} \xi, \eta, \xi \\ \xi \end{cases} = \frac{1}{8\pi} (E^2 - B^2), \end{cases} \begin{cases} \begin{cases} \frac{1}{8\pi} (\varphi^2 - A^2), \text{ (where it holds: } \xi \end{cases} \end{cases}$  $k_o = 2\pi\mu c/\hbar$ ) instead of the Lagrangian L =  $\{+\}+1$  (where 1 denotes the invariant part of the Lagrangian, which determines the interaction of the meson field with the nucleons). The

equations of the nonlinear meson field may be ascertained from the variational method  $\delta$  of div = 0. They have the form curl  $\vec{H} - \frac{1}{c} \frac{\partial \vec{D}}{\partial t} + k_o^2 A^* = 4\pi \vec{J}$ , curl  $\vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$ ,  $\vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} = \nabla \varphi$ ,

div  $\vec{D}$  +k<sup>2</sup>  $q^{*}$  =  $4\pi q$ , div  $\vec{B}$ =0,  $\vec{B}$ =curl  $\vec{A}$ ,  $\vec{D}$ = $4\pi \frac{\partial \vec{q}}{\partial \vec{E}} = \frac{\partial L}{\partial \vec{k}} \vec{E}$  +

Card 1/4

507/20-128-3-22/58 Propagation of Vector Waves in Nonlinear Mesodynamics  $+ 8\pi(\overrightarrow{E} \overrightarrow{B}) \frac{\partial L}{\partial \eta} \overrightarrow{B} + 4\pi \frac{\partial 1}{\partial \overrightarrow{E}}, \overrightarrow{H} = -4\pi \frac{\partial 0}{\partial \overrightarrow{B}} = \frac{\partial L}{\partial \xi} \overrightarrow{B} + 8\pi(\overrightarrow{E} \overrightarrow{B}) \frac{\partial L}{\partial \eta} \overrightarrow{E} -4\pi \frac{\partial l}{\partial B}$ ,  $A^* = \frac{\partial L}{\partial \xi} A$ ,  $G^* = \frac{\partial L}{\partial \xi} G$ ,  $J = \frac{\partial l}{\partial A}$ . The totality J, iq (where q denotes the density of distribution of the electrons) forms a four-dimensional current which satisfies the equation of continuity  $\frac{\partial Q}{\partial L} + \operatorname{div} \vec{j} = 0.$ The velocity of propagation of the small disturbations of the meson field may be calculated by the method of characteristics (e.g. by the method of weak unsteadinesses by Levi-Civita (Ref 3), or by the method of wave equations (Ref 4)). All these methods are adequate and deliver the same formulas for the velocity of propagation of the disturbances. The following expression is obtained from the above set of equations:  $\begin{array}{l} c\,\lambda_{\overrightarrow{H}}\times\stackrel{\rightarrow}{n}\,+\,\lambda_{\overrightarrow{D}}v\,=\,0,\;c\,\lambda_{\overrightarrow{E}}\,\times\stackrel{\rightarrow}{n}\,-\,\lambda_{\overrightarrow{B}}v\,=\,0,\;\text{where it holds:}\\ \lambda_{\overrightarrow{B}}^{\downarrow}v\,=\,-\left[\begin{array}{c} \overrightarrow{\partial B}\\ \overrightarrow{\partial t}\end{array}\right]\,,\;\lambda_{\overrightarrow{D}}^{\downarrow}v\,=\,-\left[\begin{array}{c} \overrightarrow{\partial D}\\ \overrightarrow{\partial t}\end{array}\right]\,,\;\lambda_{\overrightarrow{H}}^{\downarrow}n\,=\left[\begin{array}{c} \overrightarrow{V}\stackrel{\rightarrow}{H}\end{array}\right]\,,\;\lambda_{\overrightarrow{E}}^{\downarrow}n\,=\left[\stackrel{\rightarrow}{\Delta}\stackrel{\rightarrow}{E}\right]\,.\\ \text{The vectors }\lambda_{\overrightarrow{H}}^{\downarrow},\;\lambda_{\overrightarrow{E}}^{\downarrow},\;\lambda_{\overrightarrow{B}}^{\downarrow},\;\lambda_{\overrightarrow{D}}^{\downarrow}\;\text{are the vectors with the components} \end{array}$  $\lambda_{H_{\underline{x}}}$ ,  $\lambda_{H_{\underline{y}}}$ ,  $\lambda_{H_{\underline{y}}}$ , etc, and  $\tilde{n}$  denotes the unit vector perpendicu-Card 2/4

Propagation of Vector Waves in Nonlinear Mesodynamics 30V/20-128-3-22 inc

lar to the surface of weak unsteadiness. Further, the following relations are obtained:  $\lambda \hat{D} = \lambda \hat{E} = 0$ ,  $\lambda \hat{A} \times \hat{n} = 0$ , and

$$-\lambda_{\vec{D}} v = \begin{bmatrix} \frac{\partial}{\partial t} \left\{ \frac{\partial L}{\partial \xi} \vec{E} + 8\pi (\vec{E} \vec{B}) \frac{\partial L}{\partial \eta} \vec{B} \right\} \\ -\lambda_{\vec{H}} v = \begin{bmatrix} \frac{\partial}{\partial t} \left\{ \frac{\partial L}{\partial \xi} \vec{E} + 8\pi (\vec{E} \vec{B}) \frac{\partial L}{\partial \eta} \vec{E} \right\} \end{bmatrix}$$

The determinant of this set is equal to zero, which leads to an algebraic equation of fourth order with respect to v. The only nonlinear mesodynamics which does not exhibit double refraction or polarization of vector waves in the directions perpendicular to the collinear fields E and H has the form

$$L = \frac{C}{4\pi} \left( 1 - \sqrt{1 - \frac{8\pi F}{D}} - \frac{\eta}{c^2} + C_1, \text{ where } C \text{ and } C_1 \text{ in general} \right)$$

are functions of { . Hence, experimental investigation of the polarization of vectorial meson waves may indicate the physical reality of the electrodynamics and mesodynamics of Bonn-Infeld. If the mass and charge of the meson are known, the actual field radius and the maximum field strength may be determined in this nonlinear mesodynamics. There are 5 Soviet references.

Card 3/4

Fropagation of Vector Waves in Nonlinear Mesodynamics. SOV/20-128-3-22/58

ASSOCIATION: Moskovskiy gosudarstvennyy universitet im. M. V. Lomonosova

(Moscow State University imeni M. V. Lomonosov)

PRESENTED: May 18, 1959, by S. A. Vekshinskiy, Academician

SUBMITTED: May 18, 1959

Card 4/4

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Skobelkin, V. I.

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AUTHOR: TITLE:

Variational Principles of the Determination of the Principal Characteristics of a Ferromagnetic Phody on the Basis of the

Calculation of Its Hysteresis Loop

PERIODICAL:

Doklady Akademii nauk SSSR, 1960, Vol 150, Nr 5, pp 1012-1014

(USSR)

ABSTRACT:

The theory discussed in the present paper permits a quantitative calculation of the configuration of the whole hysteresis loop. Furthermore, it connects the properties (constants) of the ferromagnetic body with the shape of this hysteresis loop. In the case of irreversible magnetization the entropy S of a ferromagnetic body varies, so that the process of magnetization is to be described by the use of the second law of thermodynamics. The existence of magnetostrictive phenomena in a ferromagnetic body changes its surface o so that in this case an improper problem of variation (with variable limits) is obtained. Here, the field is assumed to be determined by the Cartesian coordinates Y, Z, x,  $\psi$ ,  $\varphi$  are chosen as independent coordinates. In this case, the variational equation of the

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APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

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Variational Principles of the Determination of the Principal Characteristics of a Ferromagnetic Body on the Basis of the Calculation of Its Hysteresis Loop

5/020/60/130/05/016/061 B013/B014

field reads  $\Delta(Y,Z)=\delta E^*+\int_{\mathbb{R}}\mathbb{T}\delta S d\omega =\int_{\mathbb{R}}L^*\frac{D(Y,Z)}{D(\psi,0)}\delta(\psi \theta)dx=0.$ 1 denotes the boundary of the surface  $\sigma$ ,  $L^*=w+u=$ 

 $\overrightarrow{B}(\sigma)$   $= \int_{0}^{\infty} \overrightarrow{B} d\overrightarrow{H}$  the generalized Lagrangian of the field, and it holds  $\overrightarrow{B}(\sigma)$ that  $\overrightarrow{H} = -\partial L^{*}/\partial S$ .  $\overrightarrow{B} d\overrightarrow{H}$  is equal to the magnetic pressure

acting upon the surface  $\sigma$  of the ferromagnetic body.  $\delta E^{\pi}$  denotes the complete variation in consideration of the variation in entropy and in the magnetic induction flux on the surface  $\sigma$ . The above equation is satisfied with any surface  $\sigma$  irrespective of its variation during magnetization. The variation in  $\sigma$  is to be considered by introducing an additional condition inferred from the second law of thermodynamics. The entropy variation

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e detroction is the father files to the file of the file of the interest the definitions and de-

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Variational Principles of the Determination of the Principal Characteristics of a Ferromagnetic jody on the Basis of the Calculation of Its Hysteresis Loop

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entering into the above equation must vanish for a true state. The infinitely small deviations from the true state may be regarded as possible non-equilibrium states of the ferromagnetic body. With  $\delta S=0$ , the positive magnetic flux  $\Phi^+$  through the surface  $\sigma$  has a maximum. It further holds that  $\Phi^+=\int_{-\pi}^{\pi} \frac{1}{B} \frac{1}{n} \xi d\sigma$ ,

with S=1 for  $\overline{B}$   $\overline{n}>0$  and with E=0 for  $\overline{B}$   $\overline{n}<0$ . In the general case, the thermodynamic potential of the unit of volume of a ferromagnetic body reads  $\overline{\mathcal{P}}=\mathcal{P}_0(\mathbb{N})+\mathcal{E}_{an}-H\overline{M}-\frac{H^2}{8\pi}+\mathcal{E}_{mel}$ 

magnetic body, M the magnetization of the unit of volume, magnetic body, M the magnetization of the unit of volume, an anisotropy energy. Sel elastic energy, and mel magnetical elastic energy. San is explicitly written down for uniaxial and cubic crystals. If \$\psi\$ is assumed to be a maximum, it holds

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Variational Principles of the Determination of the Principal Characteristics of a Ferromagnetic Body on the Basis of the Calculation of Its Hysteresis Loop

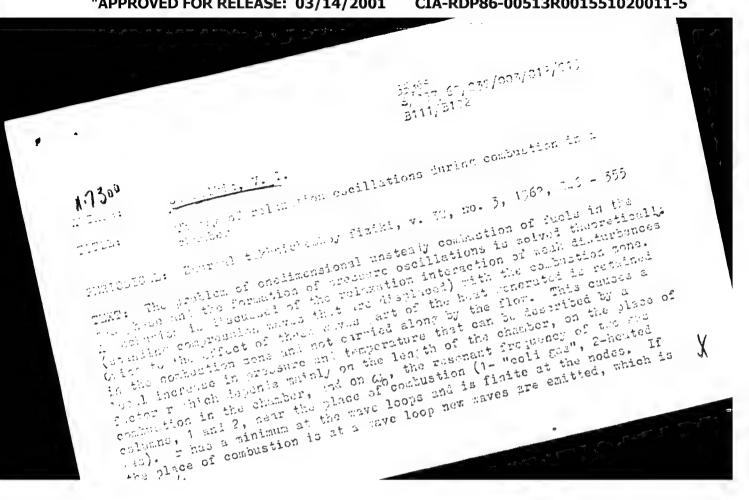
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that  $\delta = (\vec{H} + 4\pi\vec{M}) \vec{n} d\sigma^{\dagger} = 0$  with the auxiliary conditions  $\vec{u}_{ik} = (\vec{\theta} \vec{v}_{ik}) \vec{v}_{oR}$ ,  $\vec{u}_{ik}$  denotes the deformation tensor of

the ferromagnetic body. With  $\vec{N} = \vec{N}_0 + \vec{m}$  (where  $\vec{N}_0$  is the magnetic terromagnetic body. With  $\vec{N} = \vec{N}_0 + \vec{m}$  (where  $\vec{N}_0$  is the magnetic terromagnetic body. With  $\vec{N} = \vec{N}_0 + \vec{m}$  (where  $\vec{N}_0$  is the magnetic terromagnetic body. With  $\vec{N} = \vec{N}_0 + \vec{m}$  (where  $\vec{N}_0$  is the magnetic terromagnetic terromagnetic

ASSOCIATION: Moskovskiy gosudarstvennyy universitet im. M. V. Lomonosovs (Moscow State University imeni M. V. Lomonosov)

"APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5



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Thory of relacation ...

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not the case it a mode. The initial equation is a differential equation for the leviation of processes of from the equilibrium value which had been lurived by 7. I. Skobelkin (Rof3: Zorkh, 36, no. 9, 1961) (The equation only notice of the accomption that the wavelength of the disturbance is such longer than 'he width of the combustion zone). For Pi the following formum lation is given:

 $P_{1}^{t} = f_{1}\left(t + \frac{z_{1}}{a_{1}}\right) + f_{1}\left(t - \frac{x_{1}}{a_{1}}\right),$ 

where i = 1 refers to the "cold gis", i = 2 to the heated sas. The loundary conditions are:  $f_i(t) = (-1)^{Vi}F_i(t)$ , where  $V_i = 0$  corresponds to a closed charbor, and  $v_{ij}=1$  to an open one. The differential equation is solved with  $P_i = C_i^k e^{kt} \left( e^{kx/a_i} e^{a_i} + t + x_i \right) + (-1)^{i} \cos x_0 \left( -\frac{x_i}{x_i} + t + x_i \right)$ where  $C_{i}^{*}$ ,  $\chi_{i}$  are orbitrary real quantities, and a is the velocity of sound. The attenuation factor k and the frequency & satisfy complex transcendental equitions. The ownerst solution is obtained by superposition of various

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(16)

Organizers a solution of absorption in the chamber the somethist assisted expression  $P_i'=e^{-\gamma_{in}t}C_{1n}^*e^{(k_n-\gamma_{in}a_i)t}\left\{\frac{k_nz_i}{e^{-a_i}}\cos\omega_n\left(\frac{z_i}{a_i}-t+a_{in}\right)+\cdots\right\}$ 

 $\left. - (-1)^{n_i} e^{-\frac{k_n x_i}{\alpha_i}} \cos \omega_n \left( -\frac{x_i}{\alpha_i} + t - \alpha_{in} \right) \right\}.$ 

follows for  $P_1^*$ , where  $P_2^*$  is the absorption coefficient (i = 1.2). It follows from (16) that the oscillation becomes unsteady when  $k_1 - b_1 > 0$ , and that it is attenuated when  $k_1 - b_1 < 0$ . The general absorption coefficient is  $b = pa = h_1 \omega^2 + h_2 \omega^{1/2}$ , where  $h_1$  and  $h_2$  are coefficients tenending on the physical and chemical parameters of the fuel mixture.  $h_1 - h_1 = 0$  has two roots, the so-called critical frequencies  $h_1$  and  $h_2$ .  $h_1 - h_2 = 0$  has two roots, the so-called critical frequencies  $h_1$  and  $h_2$ . While the chember length  $h_1$  corresponding to  $h_1$ , the combustion is absolute with the chember length  $h_1$  corresponding to  $h_2$ , the combustion is absolute relatively steady, i. e. only for certain harmonics. For  $h_1 > 0$ , the intensificant  $h_1$  and  $h_2$  are  $h_1$  and  $h_2$  and  $h_3$  are  $h_4$ .

Thiory of relaxation ...

8/057, 32,052/005/012/019 8111/8102

TO THE RESIDENCE OF THE PROPERTY OF THE PROPER

thing the area are oscillations is relaced considerably. A conject elimination of the oscillations is, however, not attained. There are 9 references: 7 Soviet and 2 non-Soviet. The reference to the English-language publication reals as follows: ... Zucrow, J. Osborn. Jet Propulsion, 28, no. 10, 574 - 579, 1958.

ASSOCTATION: Institut Phisioheskoy fiziki AN SSSR (Institute of Chemical Physics, AS USSR)

JURNITTED: Formary CO, 1961 (initially) May 13, 1961 (after revision)

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Relaxative interaction ...

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of the minure; where rate of the reaction; z = effective width of the comboution zone; p = presoure;  $\xi$  = activation energy;  $t_r$  = relaxation

time; in a 1 refers to unbarnt, and index 2 to burnt gas. The result also halfs possible for liquid or solid fuels with combustion in the passeous above. The relaxation may cause combustion instability in rocket projuttion. The present theory refers to ideal gas, taking no account of the absorption of waves trough heat conduction, viscosity, and relaxation processes in the combustion zone. Have absorption may be important in high frequency oscillations and will be specially examined.

ASSOCIATION: Abeliening manh book, Institut khimicheskoy fiziki (Academy of Sciences USSR, Institute of Chemical Physics)

Sedwitt Dr February co, 1961

Cord J'r

SKOBELKIN, V.I.; BOLDIN, A.A.

Functions of the distribution of concentrations within the cell.

Dokl.AN SSSR 145 no.6:1396-1399 Ag '62. (MIRA 15:8)

1. Institut khimicheskoy fiziki AN SSSR. Predstavleno akademikom V.N.Kondrat'yevym. (CELLS)

SKOBELKIN, V.N., red.; BADEYAN, A., tekhn. red.

[Dairy industry of the Armenian S.S.R. during the last forty years] Molochnaia promyshlennost' Armianskoi SSR za 40 let. Erevan, Armianskaia SSR. Sovet narodnogo khoziaistva, 1961. 17 p. (MIRA 16:11)

1. Erivan. Vystavka dostizheniy narodnogo khozyaystva Armyanskoy SSR.

(Armenia-Dairy industry)

GINDINA, H.M.; KOGANOVA, G.V.; LARICHEVA, G.M.; MELKOVA, A.Ye.; POLYAKOVA, M.G.; SKOBELKINA, I.F.; IKONNIKOV, V.V., prof. otvetstvennyy red. ROSHCHINA, L., red.izd-ve; LEBEDEV, A., tekhn.red.

[State Bank of the U.S.S.R.; a brief account on the fortieth auniversary of the October Revolution] Gosudarstvennyi bank SSSR; kratkii ocherk k sorokaletiiu Oktiabria. Moskva, Gosfinizdat, 1957. 254 p. (MIRA 11:2)

 Gosudarstvennyy bank, Moscow. (Banks and banking)

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USSR/Optics - Photometry, Colorimetry, and Illumination Engineering, K-10

Abst Journal: Referat Zhur - Fizika, No 12, 1956, 35944

Author: Ryabov, I. I., Skobelev, V. M.

Institution: None

Title: New Starting-Regulating Apparatus for Luminescent Lamps

Original

Periodical: Svetotekhnika, 1956, No 1, 22-23

Abstract: None

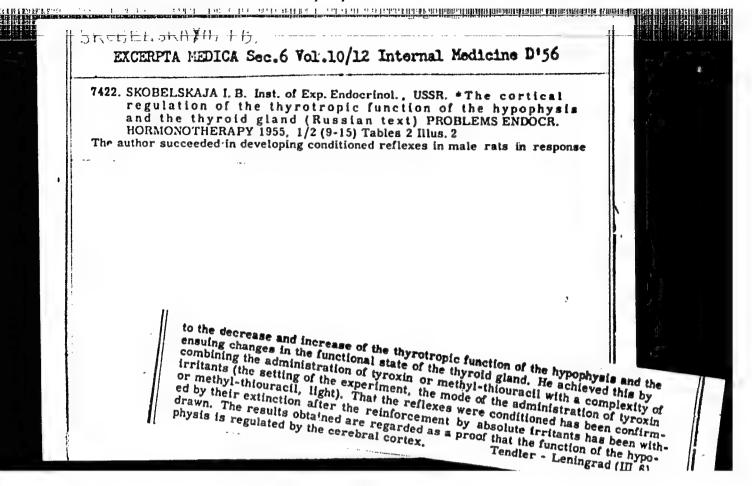
Card 1/1

TO THE RESIDENCE OF THE PROPERTY REPORTS AND THE PROPERTY OF T

BOKUCHAVA, M.A.; SKOBELEVA, N.I.

Study of volatile aldehydes of the tea plant. Dokl. AN SSSR 112 no.5:896-898 F '57. (NLRA 10:4)

1. Institut biokhimii im. A.W. Bakha Akademii nauk SSSR. Predstavleno akademikom A.I. Oparinym. (Tea) (Aldehydes)



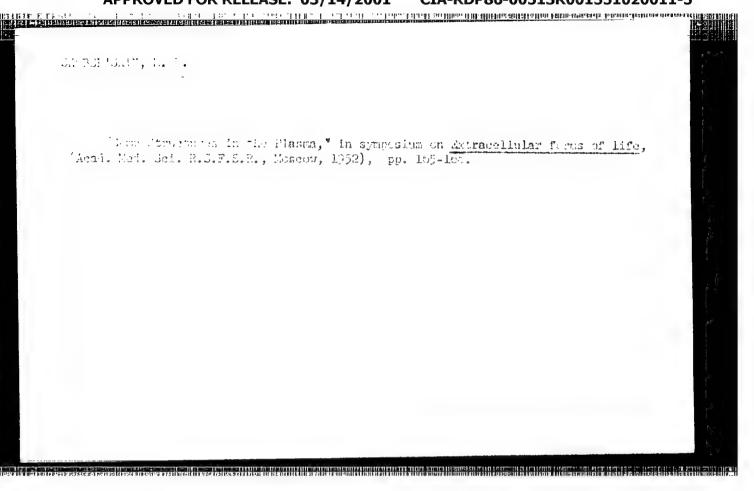
PAVLOV, Ye.; SKOBEL'SKAYA, Yu.; SAKHATSKAYA, T.

Symposium on the formation of endocrine functions in ontogeny.
Usp. sovr. biol. 60 no.2:316-319 S-0 '65. (MIRA 18:10)

SMCBELICETY, A. T.

"Comparative Study of Cytological Differentiation in the Phod Chements of Cartain
Vertetrates." Thesis for Target of Card.
Phological Sci. Sub f Jun ff, Noscow City
Pedagogical Inst imeni V. P. Potenkin

Sum ary 71, h Sep 52, Pissertations Presented for Learnes in Science and Engineering in Foscow in 1950. From Vechernyaya Moskva,
Jen-Pec 1950.



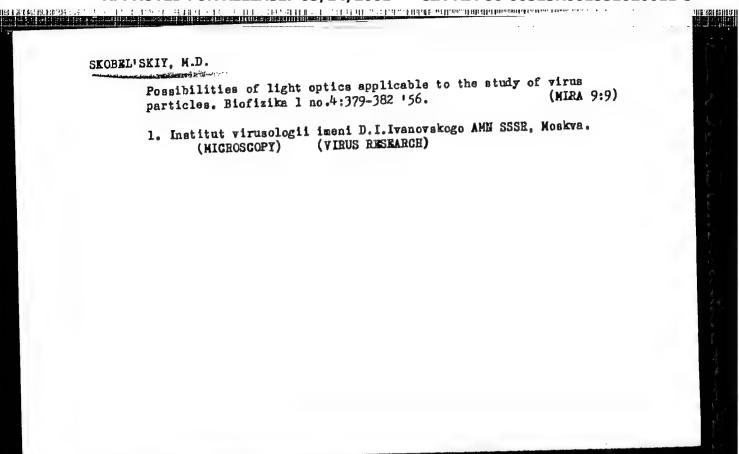
#### SKOBEL'SKIY, M.D.

Development of vaccinia virus outside of the organism. Biul. eksp. biol. i med. 38 no.11:66-70 N '54. (MLRA 8:1)

1. Iz Instituta virusologii imeni D.I.Ivanovskogo (dir. deystvitelinyy chlen AMN SSSR prof. M.P.Chumakov) AMN SSSR, Moskva.

(VACCINIA, virus, develop. outside of organism) (VIRUSES,

vaccinia virus, develop. outside of organism)



的<mark>技术技术和连接的技术</mark>的技术的大学,一直不是这个企业,不是有关的企业,就是有关的企业,但是一个工程的自己的工作,但是一个工程的自己的工作,但是一个工程的工作的工作,并不 E-3 USSR/Virology - Human and Animal Viruses : Referat Zhurn - Biol. No 16, 25 Aug 1957, 68242 Abs Jour : Skobelskiy, M.D., Avakyan, A.A., Burak, A.I. : Use of the Luminescent Method of Microscopy for Perfecting of Laboratory Diagnosis of Poliomyelitis. Author Title : Vopr. Virusologiy, 1956, No 6, 58. Orig Pub : Hela(?) cells were cultured on cover glasses, filtrated, treated with acridine yellow (1:1000) and studied through a luminescent microscope. Regardless of the age Abstract of the culture, the protoplasm luminesced by a dull-green light, surrounding the far more luminous nucleus and small nuclei. Six hours after infection of the poliomyelite by the virus, thin sprouts of cells began to manifest orange-yellow illumination, which further enveloped adjacent portions of the protoplasm spreading from the periphery of the cell toward the nucleus. On the second day after the infection of the protoplasm, the nucleus - 10 -Card 1/2

USSR/Virology - Human and Animal Viruses

E-3

Abs Jour

: Referat Zhurn - Biol. No 16, 25 Aug 1957, 68242

also began to luminesce with a flaming-yellow light. This luminescence did not arise when the specific immunizing serum was introduced into the culture together with the virus.

Card 2/2

- 11 -

SKOBEL'SKIY, M.D.; BURAK, A.I.

Method of arranging tissue culture for morphological analysis.

Vop. virus 5 no.4:494-496 Je-Ag '60. (MIRA 14:1)

l. Institut virusologii imeni D.I.Ivanovskogo AMN SSSR, Moskva. (TISSUE CULTURE)

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001551020011-5"

KHONDKARIAN, O.A.; SKOBEL'SKIY, M.D.; KHVAN, L.M.; BURAK, A.I.

Clinical aspects and etiology of acute serous meningitis. Vest. AMN SSSR 17 no.7:13-17 '62. (MIRA 15:10)

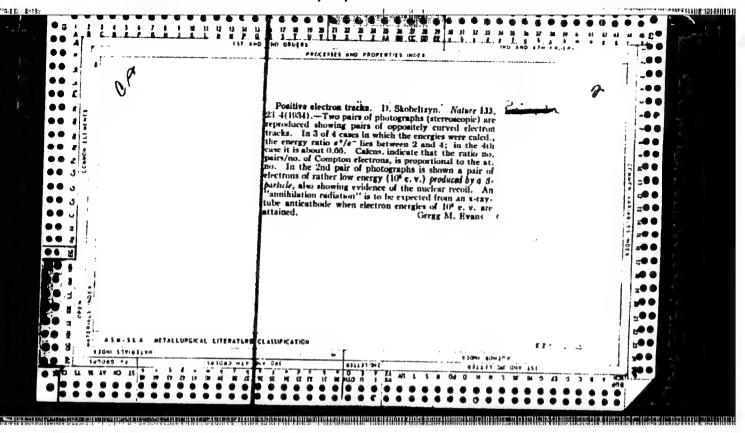
1. Institut nevrologii AMN SSSR.
(MENINGITIS) (VIRUS DISEASES)

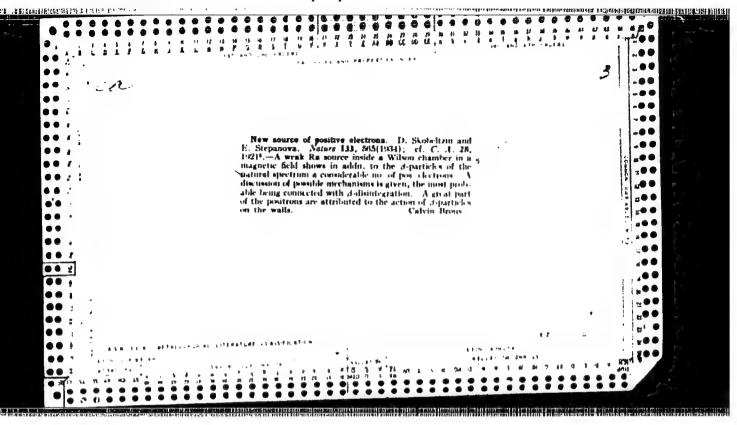
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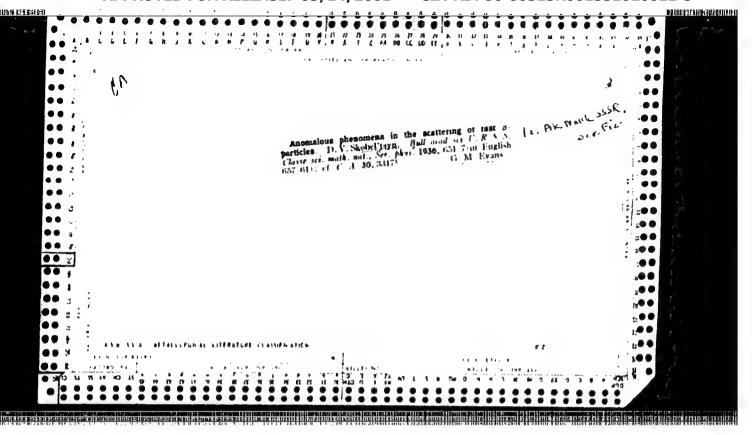
MITKALINNYY, V.I., kand.tekhn.nauk; MOLCHAROV, N.G., kand.tekhm.nauk;
Prinimali uchastiye: MEVEDOMSKAYA, I.N.; SHKOL'HIKOV, Yu.M.;
VOLVENKIN, V.K.; RAYSKIY, R.M.; BELEN'KIY, A.M.; SKOBEL'TSIR,
S.S.; FEY GZHU-HIN; CHAHAO TIN'-YUAN'

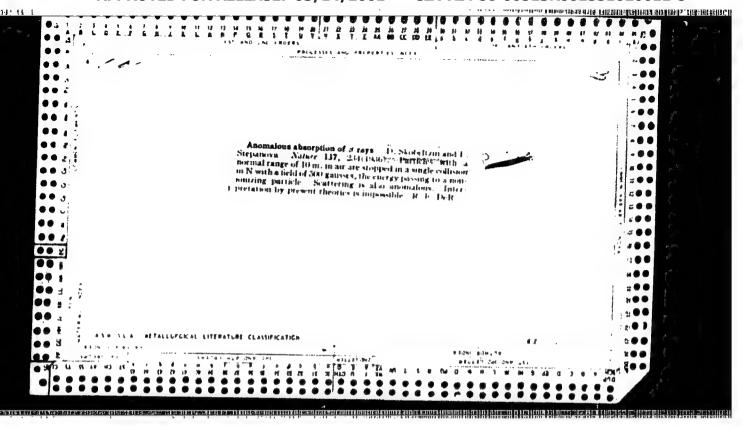
Improvement of bell-type furnaces for bright annealing. Stal'
(MIRA 15:5)

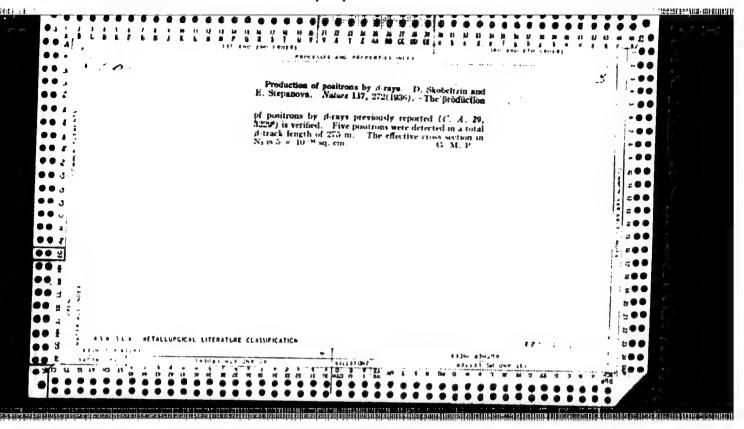
1. Moskovskiy institut stali.
(Furnaces, Heat-treating) (Annealing of metals)

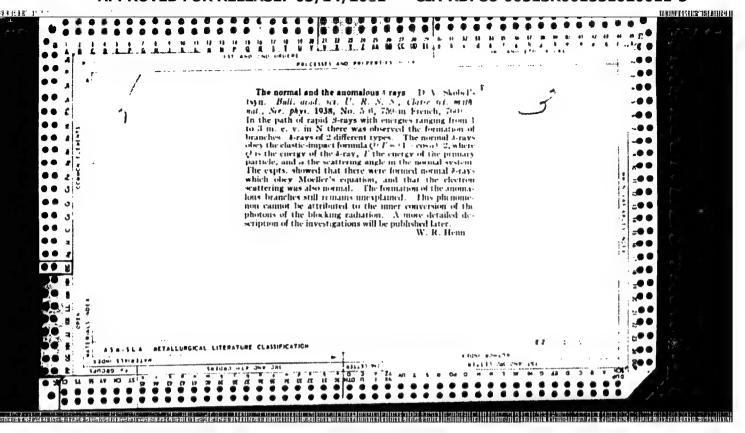


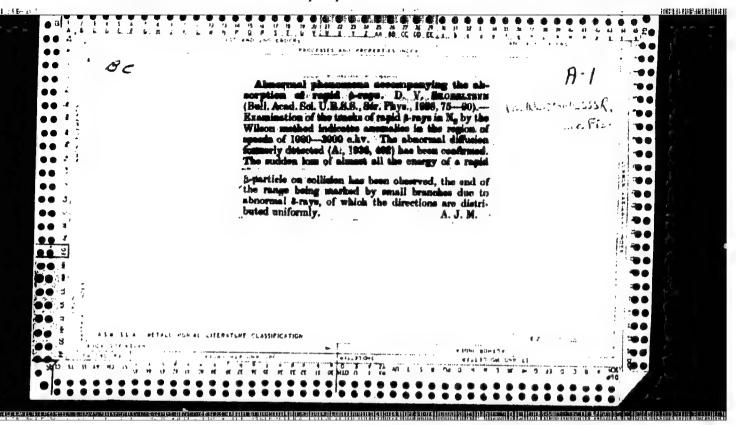


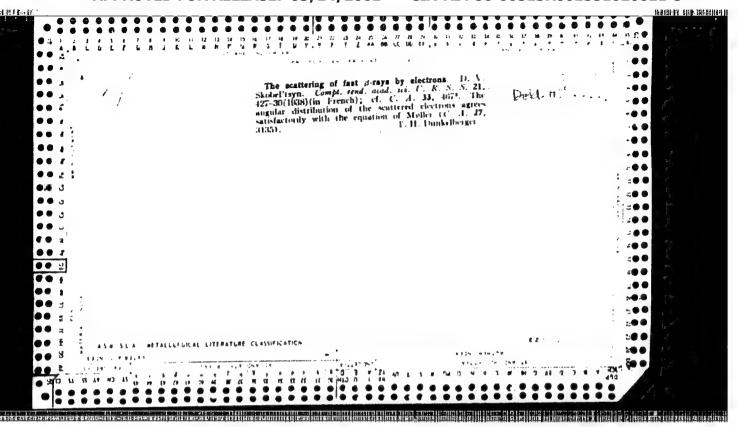


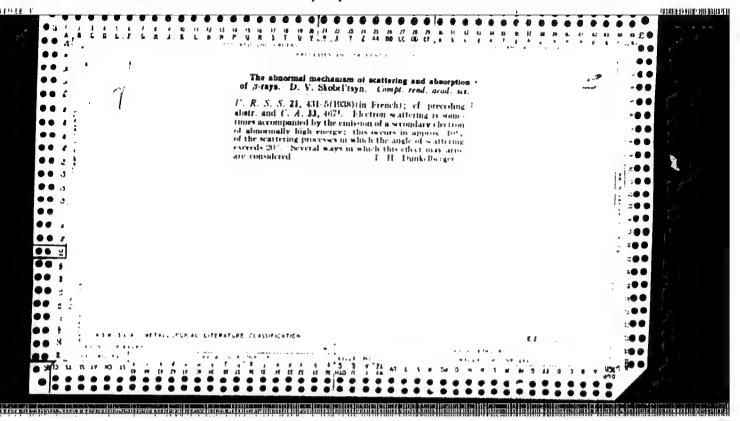


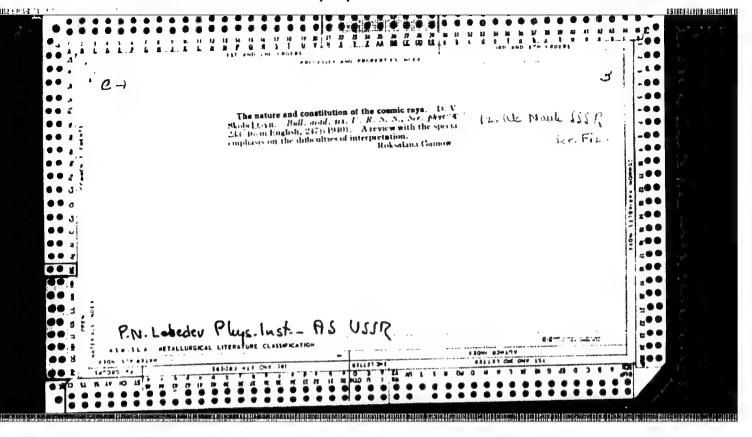


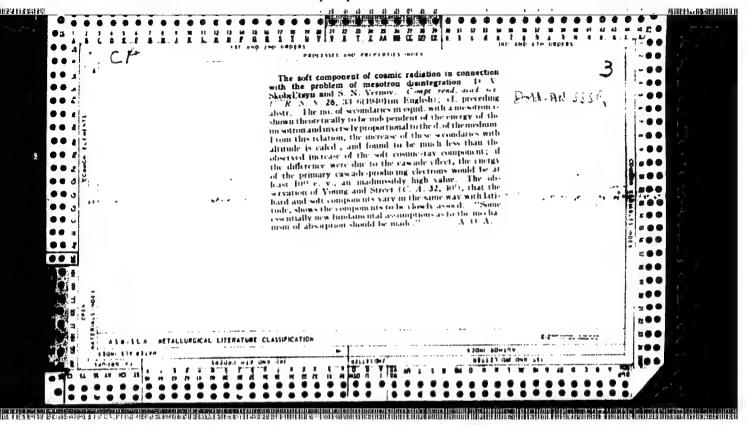


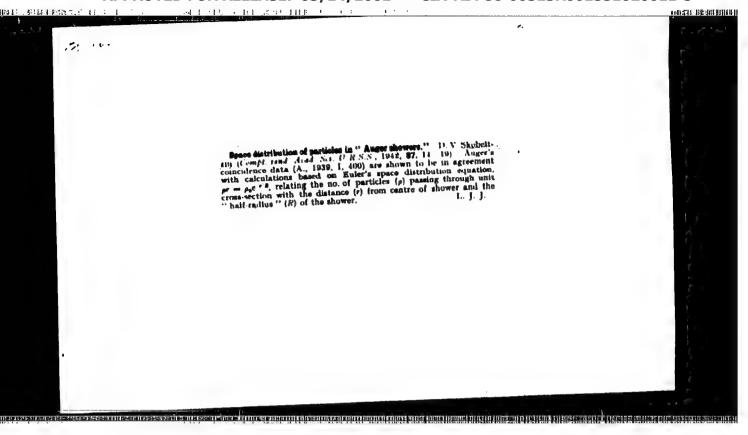


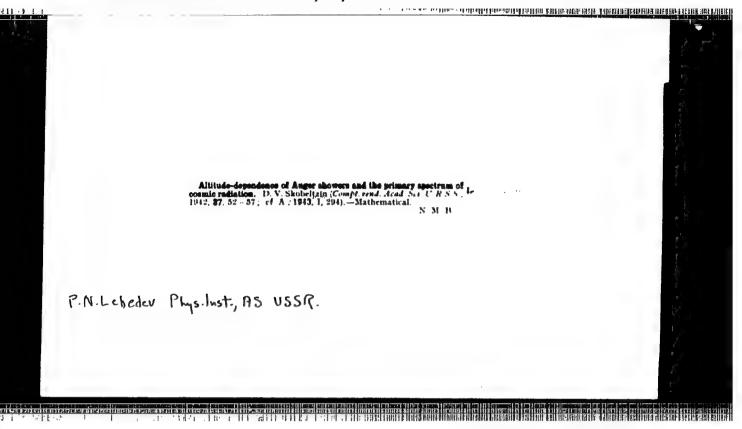


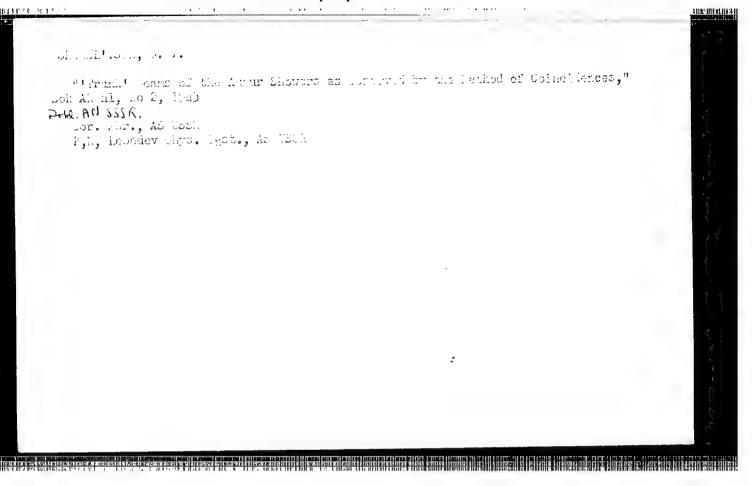


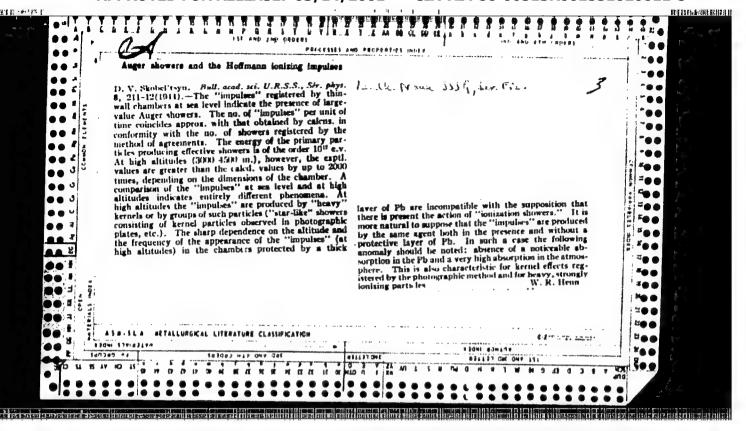


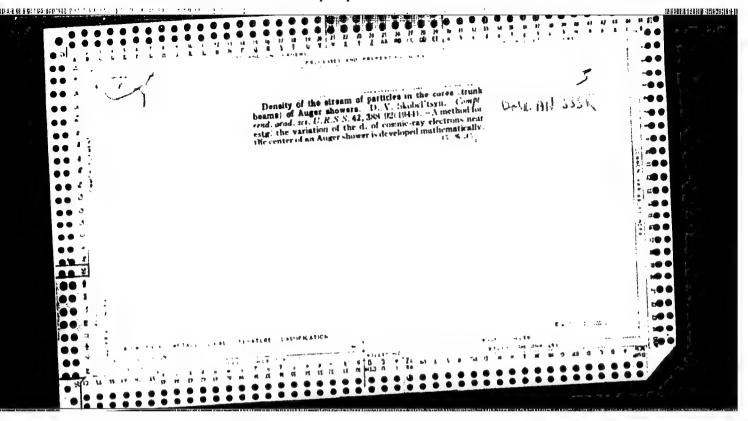


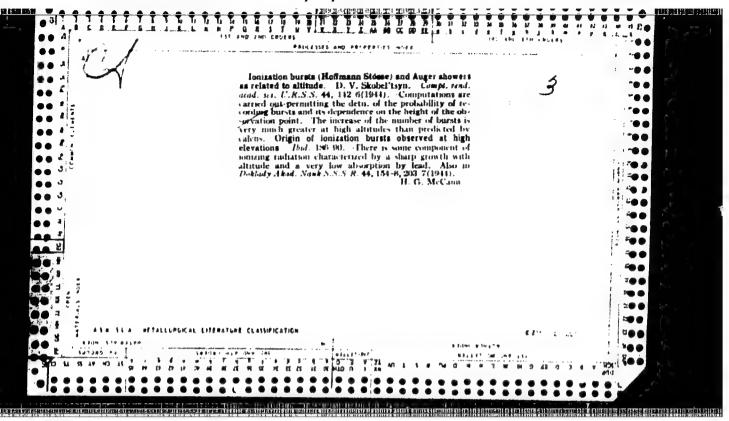


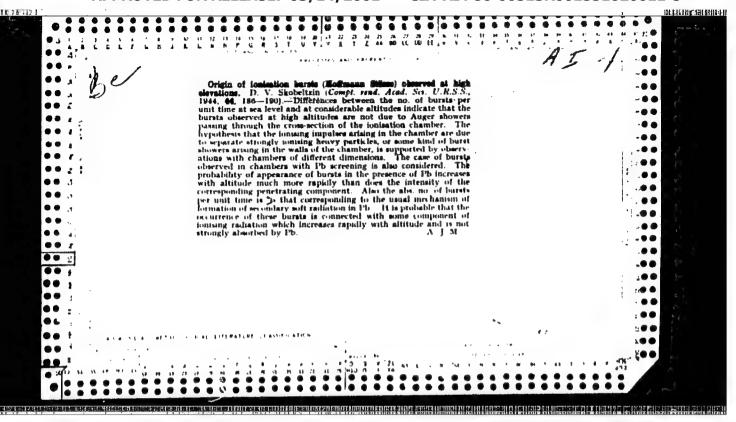


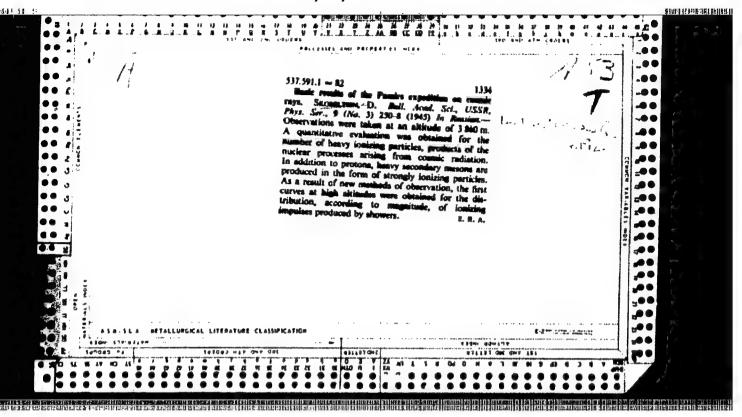


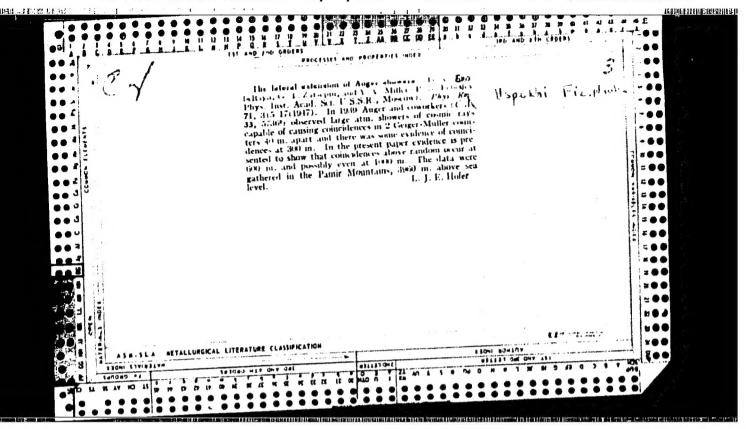












"The Atômic Comb," collaboration intelligedity, fol. 3, and edition, g. 627, 1019

GURO, G. NIKOLAEV, V... RAZORENOV, L., CHIVILO, I. and SKOVELLEVA D. V.

Altitute travel and curves of absorption in the angle of particles which generate impacts in the ionization chamber at heights os 3860 and 4700 meters. (Presented by Academician D.V. Skobeltsyn\* 20 May 1949.)

Reports of the academy of Sciences USSR Vol. 57, No.3, Sept. 23, 1949.

The state of the s					men eine Anadetinea Jiobülles 62, 711 (2005)
J. TELITEYN, D. V.	successive application of the avalanche exceeded experimental results only by a Submitted 12 May 49.	USSR/Nuclear Physics (Contd)	Observations at 5,860 meters showed dences caused by atmospheric showers are could be observed in separating distance of about 1,000 meters than limits. When such coincidences were	"Critically Extensive Atmospheric Showers P. Radiation," Acad D. V. Skobel'tsyn, Phys P. N. Lebedev, Acad Sci USSR, 4 pp	USSR/Nuclear Physics Commic Rays Geiger Counters
Z6164/45	theory, effect small percent.	54/49 <u>195</u> Jul 49	in Geiger count- the counters by a within previous calculated by	s of Cosmic	64 Lak